

Equilibrium and folding simulations of NS4B H2 in pure water and water/2,2,2-trifluoroethanol mixed solvent: examination of solvation models

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Received: 8 April 2013 / Accepted: 23 June 2013 / Published online: 7 July 2013
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Abstract The structural stability and preference of a protein are highly sensitive to the environment accommodating it. In this work, the solvation effect on the structure and folding dynamics of a small peptide, NS4B H2, was studied by computer simulation. The native structure of NS4B H2 was solved previously in 50 % v/v water/2,2,2-trifluoroethanol (TFE) mixed solvent. In this work, both pure water and water/TFE cosolvent were utilized. The force field parameters for water were taken from the TIP3P water model, and those for TFE were generated following the routine of the general AMBER force field (GAFF). The simulated structure of NS4B H2 in the mixed solvent is quite in line with experimental data, while in pure water it undergoes a large structural deformation. The generalized Born (GB) model was also investigated by tuning the dielectric constant to match experimental measurements. However, the results show that its performance was less satisfactory. Two independent direct folding simulations of NS4B H2 in explicit water/TFE cosolvent were carried out, both of which resulted in successful folding. Investigation of the distribution of solvent molecules around the peptide indicates that folding is triggered by the aggregation of TFE on the peptide surface.

Keywords 2,2,2-Trifluoroethanol · Molecular dynamics · Protein folding · Solvent model · Aggregation

Introduction

Solvents play an important role in modulating the stability of secondary structures of proteins and peptides. 2,2,2-Trifluoroethanol (TFE) is well known as one of the most efficient inducers of secondary structure, especially helical structure, in peptides. In the last few decades, the impact of fluorinated solvents on the structure of peptides and proteins has been studied extensively by both experimental [1–4] and computational [5–10] means. Various mechanisms by which TFE promotes the stabilization of particular secondary structure elements have been proposed. The dielectric constant of aqueous TFE solution is much lower than that of pure water under ambient conditions, [4] which weakens hydrophobic interactions but enhances electrostatic interactions such as hydrogen bonds [11, 12]. Therefore, an ordered structure with main chain hydrogen bonds is preferred in TFE. In addition, the aggregation of TFE molecules around the peptide gives rise to a matrix that prevents the intrusion of water molecules to the surface of the peptide [4, 6, 8]. It is well known that water is a competitive hydrogen-bond participant, which may destabilize intra-protein hydrogen bonds. Besides, the helix triggering feature may also arise from the reduced entropy penalty associated with α -helix formation after the association of TFE molecules around the protein surface [13].

An accurate description of the interactions between different components in aqueous solution and that between the protein and solvent is essential for realistic biomolecular modeling. Unfortunately, the great majority of force fields are designed for pure solutions, especially pure water solvent under ambient conditions. Several force field parameters for TFE have been proposed, for example by van Buuren and Berendsen [14, 15], De Loof [16], and Fioroni [17]. Physical

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properties calculated from these force fields are in qualitative agreement with experimental measurements. So far, none of the currently available force fields outperforms the others to well describe all the properties of pure liquid TFE and water/TFE mixture. Usually, the choice of the most appropriate force field depends on the major properties of interest. In this work, we developed a new force field for a water/TFE mixture based on the philosophy of the general amber force field (GAFF) [18], and studied its performance in modeling the structure and dynamics of a short peptide in this mixed solvent. GAFF is an indispensable tool in the AMBER community for the modeling of small organic molecules. However, its availability for modeling an organic liquid has not been examined fully. In our previous work [19], a detailed study of the performance of GAFF in modeling pure TFE liquid showed that GAFF was satisfactory for the modeling of TFE liquid, although there was still room for improvement.

Nonstructural protein 4B (NS4B) plays a critical role in the formation of the hepatitis C virus (HCV) replication complex—a relatively poorly characterized 27-kDa integral endoplasmic reticulum (ER) membrane protein [20]. It is predicted to contain a cytosolic N-terminus, two amphipathic helices extending from amino acids (AA) 6 to 29 [21] and AA 42 to 66 [22], a central portion that contains four transmembrane segments [23], and a cytosolic C-terminal part that includes two α -helices (AA 201 to 213 N and AA 228 to 254 [H2])[24]. A partial translocation of the N-terminus of NS4B into the ER lumen has been demonstrated, and this transient topology of NS4B might contribute to the induction of membranous vesicles [22, 25]. In addition, the C-terminal palmitoylation on residues Cys257 and Cys261 [26] is important for the oligomerization, which could contribute to membrane web formation and HCV RNA replication [27]. Here, the peptide studied is H2 in the C-terminal portion of NS4B, which is an amphipathic α -helix extending from Ser227 to Ser254. In the following text, this peptide sequence is re-numbered, starting from 1 for Ser27.

In this work, we compared various solvation models for a 50 % (v/v) water/TFE mixture as solvent through MD simulations of NS4B H2. First, we carried out equilibrium MD simulations of this peptide in explicit 50 % (v/v) water/TFE mixed solvent, and compared with simulations in pure water. Both explicit and implicit generalized Born (GB) solvent models were considered. The peptide structure was well maintained in explicit water/TFE cosolvent. In explicit solvent models, the peptide readily formed an ordered helical conformation in the presence of TFE. However, structural deformation was observed in pure water. This result was consistent with previous studies by other groups [6–8]. With the GB water model, the conformation of this peptide was unstable, which was in good agreement with the result in explicit water. However, it was also unstable in GB water/TFE solvent, with

a relative dielectric constant set to 53.0, compared to the result from the simulation in explicit water/TFE cosolvent. Considering that the aggregation of TFE molecules around the peptide leads to a much higher TFE concentration in the vicinity of the peptide than the nominal value of the bulk solution, the actual screening effect experienced by this peptide is not close to the macroscopic measurement. Therefore we also carried out a simulation of this peptide in GB with external relative dielectric constant set to that of pure TFE (27.1). However, the result did not improve too much. Therefore, the subsequent folding simulations were carried out only in an explicit water/TFE mixture. We drew consistent conclusions from two independent trajectories that the peptide folded in an explicit 50 % (v/v) water/TFE mixture and the final structures of both trajectories were in good agreement with its NMR structure [28].

Methods

The parameters of the TFE molecule were generated following the standard procedure of GAFF. The initial structure of TFE was optimized at HF/6-31G** level, and the atomic charges were fitted to the electrostatic potential calculated at B3LYP/cc-pvtz level with restraints applied to avoid a large amplitude of atomic charges [29]. Bonded and van der Waals parameters were determined by atom types. The 50 % (v/v) mixed water/TFE solution was generated by putting 1:4 TFE and water molecules in a periodic box utilizing the LEaP module in AmberTools. A 32 ns molecular dynamics (MD) simulation at 300 K in NPT ensemble was carried out to relax this mixture. The final structure with a density of 1.15 g/cm³ serves as the building blocks of solvent for dissolving the peptide.

The initial structure of the NS4B H2 peptide in the equilibrium simulations was obtained from Protein Data Bank (entry 2KDR [29]). AMBER03 force field parameters were used for the peptide. For simulations in explicit solvent models, the peptide was placed in a periodic truncated octahedron box of either TIP3P water [30] or a mixture of TFE and water. The minimal distance between the peptide and the water box was 1.2 nm and that between the peptide and the water/TFE box was 1.6 nm. A larger buffer size for water/TFE was necessary, because many TFE molecules were removed automatically when dissolving the peptide. Some water molecules must be removed manually to attain the required TFE:water ratio. Therefore there were some bubbles in the solution initially, and the size of the box would shrink during the simulation under constant pressure. Finally, the mixed solvent contains 4,528 TIP3P water molecules and 1,132 TFE molecules. Two chloride ions were added to neutralize each of the systems. For GB solvent models, the GB solvation model proposed by Onufriev, Bashford and Case [31] was employed. The nonpolar

solvation term was excluded. Nonbonded interactions were fully counted without any truncations. A unit dielectric constant was assigned to the solute interior, while the continuum solvent region was assigned a dielectric constant of 78.5, 53.0 and 27.1 for water, water/TFE mixture and pure TFE respectively [4].

For simulations in explicit solvent models, each system was energy minimized by a steepest descent algorithm for 5,000 steps with strong restraints imposed on the solute and was then subjected to a thorough relaxation without any restraints. The optimized structures were heated up to 300 K in 100 ps in water or 5 ns in water/TFE mixed solvent with weak restraints applied to the peptide, followed by a 300 ns production simulation at NTP ensemble. All simulations were carried out under ambient conditions by employing the Berendsen algorithm [32] to regulate the temperature and pressure with $\tau_T=0.25$ ps and $\tau_P=2.0$ ps. The SHAKE algorithm [33] was utilized to constrain all covalent bonds involving hydrogen atoms, and the integral time step of propagation was 2 fs. The particle mesh Ewald algorithm [34] was utilized for long-range interactions with a 10-Å cutoff in real space. The van der Waals interaction was truncated also at 10 Å. For simulations employing the GB model, the simulation setup was quite similar to that in the explicit water box except that only a single energy minimization procedure was applied, and the production simulation extended only to 100 ns. The folding simulation of the peptide started from a linear structure generated by the LEaP module of AmberTools, and it followed the same procedure and conditions as in the equilibrium simulation of the peptide in an explicit water/TFE mixture, except that the simulations extended to 1 μ s each. All simulations were performed using the AMBER 11 package [35].

Results and discussion

Equilibrium simulation analysis

Explicit solvent models

Shown in Fig. 1a–f are snapshots extracted evenly from the trajectory in explicit water with a 50 ns interval. The snapshot with the largest root mean-square deviation (RMSD) from the

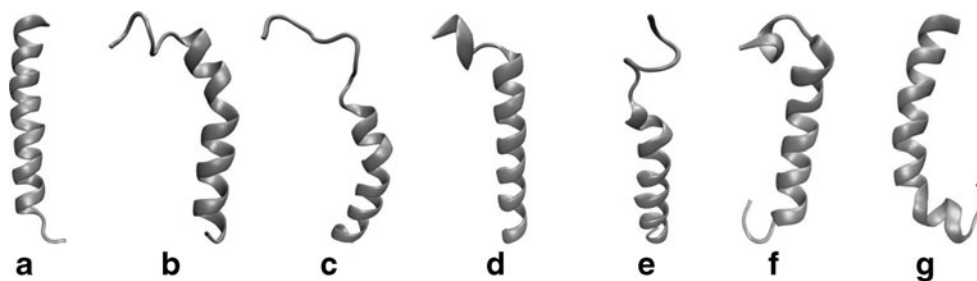


Fig. 1 Snapshots of the peptide taken every 50 ns along the equilibrium trajectory in explicit water (a–f) and the configuration with maximum root mean-square deviation (RMSD) from the first model of NMR structure in explicit water/2,2,2-trifluoroethanol (TFE) mixture (g)

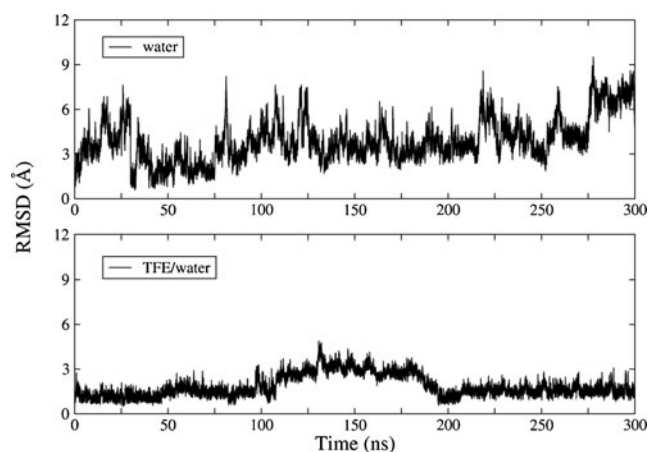


Fig. 2 Backbone RMSD from the first model of NMR structure during equilibrium simulations in (top) explicit water and (bottom) water/TFE

first model of the NMR structure in the trajectory in water/TFE mixed solvent is also shown for comparison (Fig. 1g). In explicit water, the peptide was not stable and unfolding took place at the N-terminus. At 100 ns, two of seven helical turns became random coil and the remaining five turns were distorted and bore a bent shape. Thereafter, refolding and unfolding can be observed several times, according to the time evolution of RMSD shown in Fig. 2. The peptide underwent considerable deformation with RMSD fluctuating back and forth in large amplitude. Refolding (RMSD < 2 Å) was found in the second 25 ns. In the last 25 ns, the mean RMSD was around 7 Å and the largest RMSD reached 9.6 Å. While in 50 % (v/v) water/TFE mixed solvent, the peptide was much more stable. Even for the snapshot with the largest RMSD, only a small distortion occurred to the original straight long helix. In majority of the simulation time in water/TFE mixed solvent, the RMSD was below 2 Å, which was apparently in the folded state. Only in the second 100 ns was the peptide partially unfolded, with RMSD approaching 5 Å. It refolded and stayed in the native structure in the final 100 ns. This result was consistent with experimental results in that only in water/TFE cosolvent could the structure of this peptide be determined.

The unfolding of secondary structure is accompanied by breaking of the main chain hydrogen bonds. We plot in Fig. 3 the variations in bond lengths of the four hydrogen bonds at each end of this peptide. The bond length was measured as

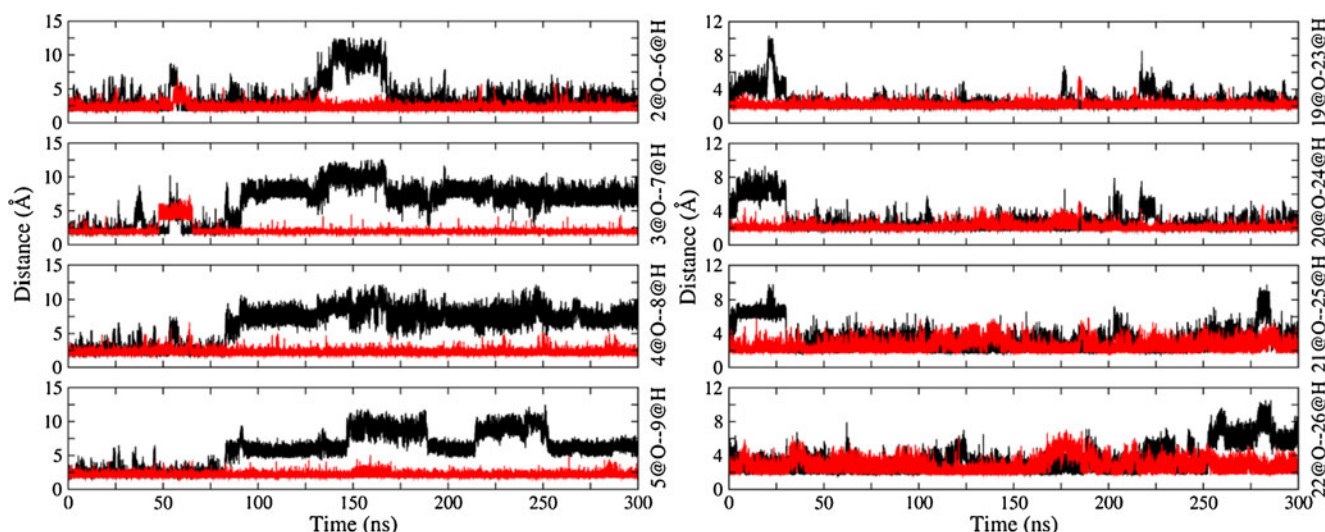
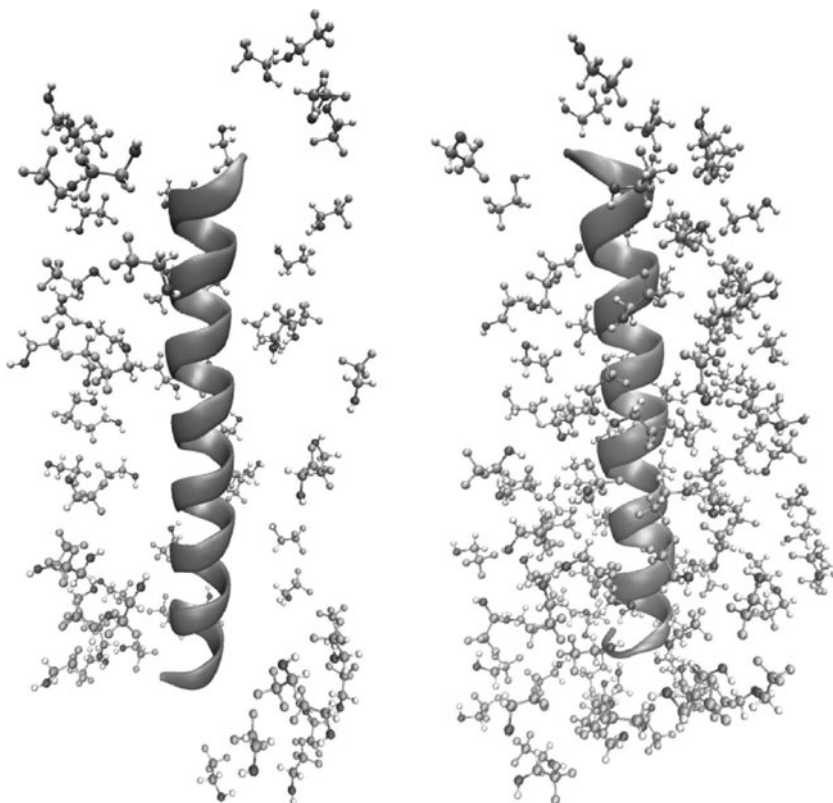


Fig. 3 Variation in hydrogen bond length in the peptide during the equilibrium simulation in explicit water (*black*) and water/TFE (*red*). Atoms are labeled as “residue number@atom name”

the distance between the hydrogen atom involved and the acceptor oxygen atom. When hydrogen-bonded, the distance between these two atoms should be below 2.5 Å. In the first 25 ns of the simulation in pure water, 3 (Leu¹⁹–Leu²³, Leu²⁰–His²⁴ and Arg²¹–Gln²⁵) of the last four hydrogen bonds were broken, and the distance between oxygen and hydrogen atoms went up to over 4 Å. However, they all refolded back in around 30 ns, and remained very stable in the subsequent simulation time. The last hydrogen bond was

very stable for 250 ns, but was broken in the last 50 ns. In general, the last four hydrogen bonds were quite stable and remained bonded for most of the simulation time. Nevertheless, the first four hydrogen bonds were vulnerable. The second to the fourth hydrogen bonds (Ala³–Val⁷, Ala⁴–Thr⁸, Ala⁵–Ala⁹) broke at about 80 ns, and never reformed. This was consistent with the observation that the N-terminus of this peptide unfolded. On the contrary, all the hydrogen bonds were very stable during the simulation in water/TFE

Fig. 4 Comparison of the first and the last snapshots in equilibrium simulation in explicit water/TFE mixed solvent. TFE molecules within 0.6 nm from the peptide are shown



mixed solvent. Only in some minor portion of time was breaking of these hydrogen bonds noted.

Therefore, TFE molecules are essential for the structural stability of this peptide. Understanding the microscopic mechanism of TFE in protecting the helical structure is valuable. In the beginning of the simulation in water/TFE mixed solvent, the water molecules and TFE molecules were distributed evenly around the peptide and in the bulk. With the simulation ongoing, TFE molecules gradually aggregated on the surface of this peptide and pushed water molecules away (see Fig. 4). This phenomenon has also been observed in some other studies [4, 6, 8], and can be explained by the amphipathic nature of TFE. The CF_3 group renders a significant hydrophobicity to the molecule and the electronegativity of fluorine increases the acidity of the hydroxyl group, which makes TFE a better hydrogen bond donor but a poorer hydrogen bond acceptor than a water molecule [36], and allows TFE to reside on the peptide surface. Therefore, the protecting mechanism of water/TFE cosolvent on the helical structure is due not only to the low dielectric constant of this solvent, which may enhance the Coulomb interaction, but also the coating effect of the TFE molecules, which blocks the invasion of water molecules to the main chain hydrogen bonds—the main cause of the destabilization of helical peptide in pure water. The local TFE concentration (LTC) around the peptide residues was calculated by extracting the relative numbers of TFE and water molecules within a 0.6 nm shell surrounding each residue and considering the average excluded volume of 0.019 and 0.07 L/mol for water and TFE, respectively [7, 37]. The LTC for each residue in the system is depicted in Fig. 5. The peak of LTC appears in the center of this peptide, which approaches 0.9. The LTC decays towards both terminals. This is reasonable because the terminal residues are usually more flexible than those in the center of a helix. The high concentration of TFE molecules protects the intra-protein hydrogen bonds from the competition of water

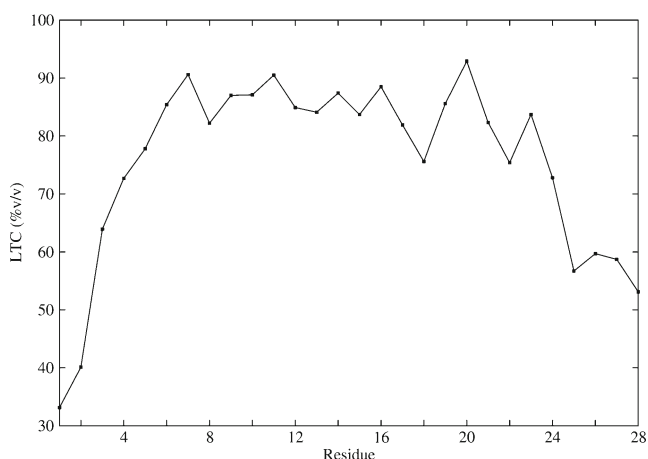


Fig. 5 Average local TFE concentration within a distance of 0.6 nm from the C_α atom of each residue of the peptide in equilibrium simulation

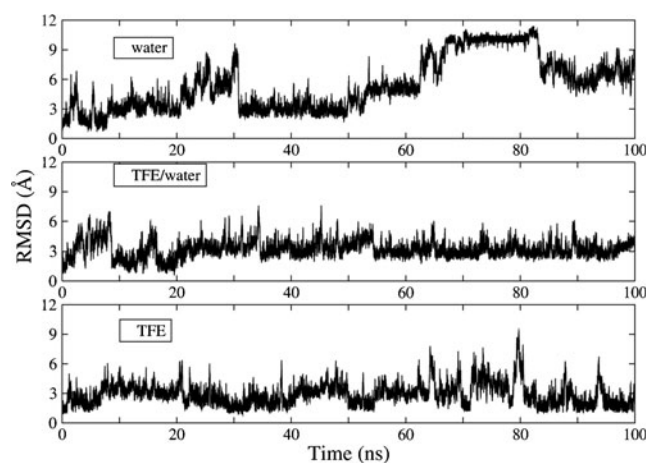


Fig. 6 Backbone RMSDs from the first model of NMR structure along the generalized Born (GB) equilibrium simulations in (top) water, (middle) water/TFE, and (bottom) TFE

molecules in forming hydrogen bonds with the protein. Therefore, the peptide maintains a helical structure.

GB solvent models

The GB solvent model, as an alternative to the explicit representation of solvent, is particularly attractive because of its algorithmic simplicity and computational efficiency. Unfortunately, the approximations in the GB model also limit its applications in molecular modeling. Therefore it is necessary to study the performance of the GB model in modeling the solvent. The GB solvation model proposed by Onufriev, Bashford and Case [31] as implemented in the AMBER package was employed, with an external relative dielectric constant ϵ set to 78.5 (water), 53.0 (water/TFE) or 27.1 (pure TFE) [4]. The time series of the backbone RMSDs from the first model of NMR structure in three simulations is shown in Fig. 6. During the simulation in implicit water, the peptide was very flexible, and it quickly unfolded, with

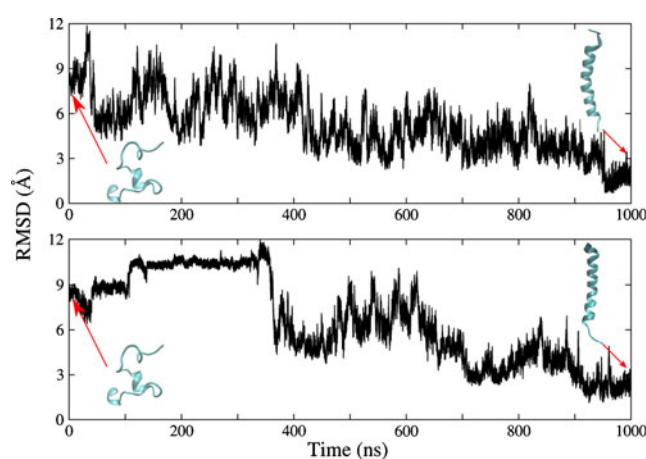


Fig. 7 Backbone RMSDs from the first model of NMR structure in two folding simulations in explicit water/TFE mixture

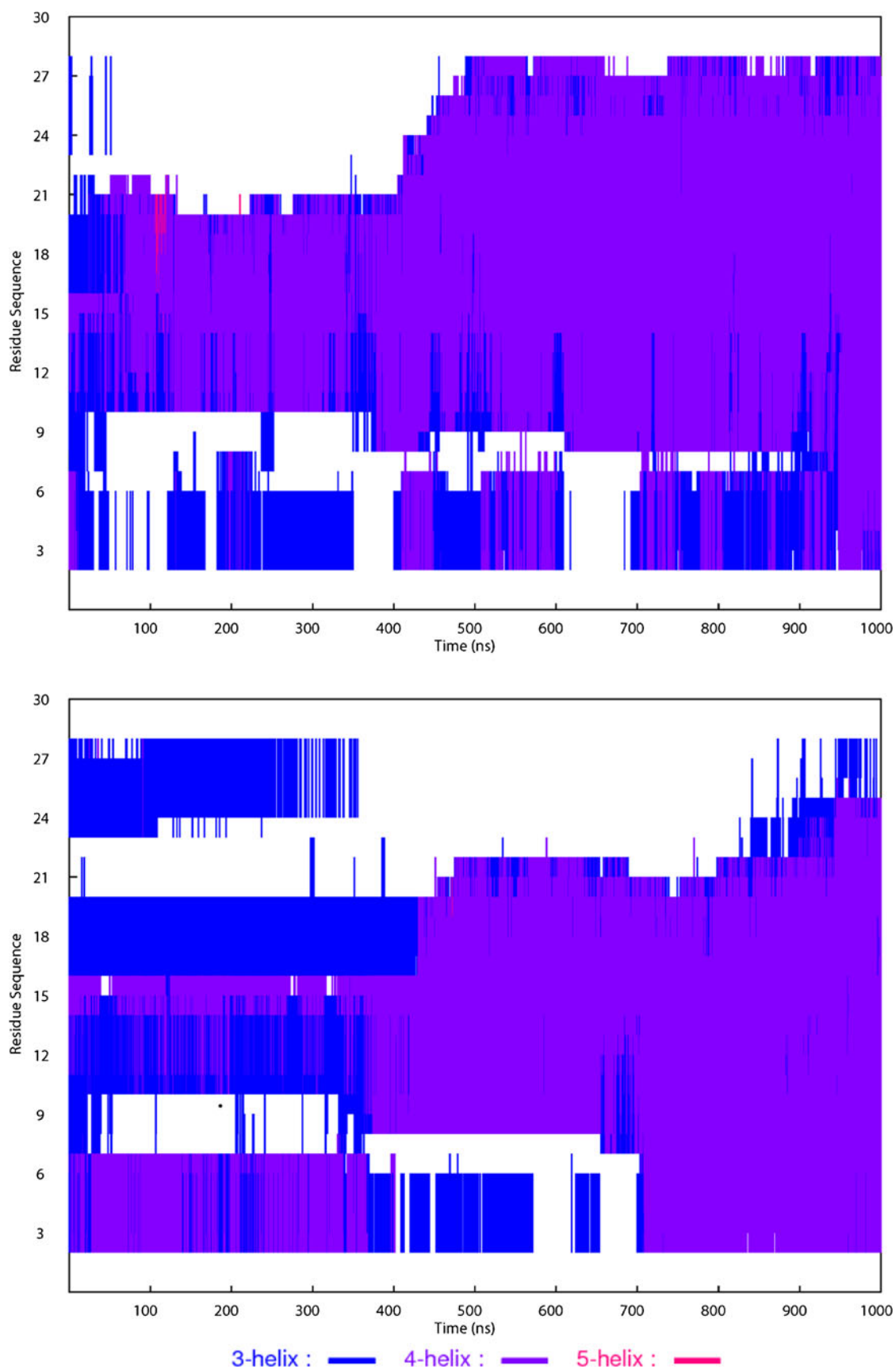


Fig. 8 Variations of the secondary structure in the folding simulations

RMSD approaching 10 Å in 30 ns. Although it also folded back to the native state very quickly, another unfolding event was observed after 50 ns, with RMSD gradually going up to over 10 Å and staying there for 20 ns. This is consistent with the simulation in explicit water, but contrasts with the much suppressed fluctuation of RMSD along the simulation in water/TFE mixed solvent. A discrepancy between simulations in GB and in explicit water/TFE mixture can still be noted. The RMSD fluctuates mainly between 2 and 4 Å, and a larger RMSD over 7 Å can also be seen. It might be argued that the first solvation shell of the peptide in mixed water/TFE is occupied mainly by TFE molecules. Therefore, the dielectric constant is different from the macroscopic measurement. To examine this conjecture, a simulation in implicit solvent with the dielectric constant set to that of pure TFE was also performed. However, the result did not show any improved stability of this peptide over that in $\epsilon=53.0$. This result shows that an implicit solvent model like GB is not a good model for TFE or water/TFE mixed solvent, and its failure is not due merely to its incapability of representing the inhomogeneous distribution of TFE and water molecules in the first solvent shell and in the bulk.

The discrepancy between explicit solvent model and GB model attributes to the limitations of GB models. The GB model is based on replacing discrete solvent molecules by an infinite continuum with the dielectric properties of the solvent, which is extracted from the fundamental discrete model by several layers of approximations, each of them adding its own limitations on the model. The algorithmic simplicity and computational efficiency, combined with relatively reasonable accuracy of the GB approximation make it attractive in many practical applications of MD simulations of biomolecules, e.g., protein folding and design [38–41], large-scale motions in macromolecules [42–44], peptides and proteins in membrane environments [45–47], and pKa prediction and constant pH simulations [48, 49] etc. However, the GB methodology in modeling more complex environments such as a water/TFE mixture requires a heterogeneous description with a spatially varying dielectric constant. This remains a challenge to the computational biology community. In addition, the dielectric constant changes during the MD simulation with the TFE molecules aggregating around the peptide, which also adds to the difficulty of the GB model.

Folding simulation in explicit water/TFE solvent

Based on the result shown above, the explicit water/TFE mixed solvent model is indispensable to the simulation. Therefore, in the subsequent folding simulations of NS4B H2 in water/TFE cosolvent, the solvent molecules were represented explicitly. Two independent simulations with different initial velocities were carried out, the results of which were quite consistent with only marginal dissimilarities. The

RMSDs of backbone atoms from the first model of the NMR structure are shown in Fig. 7. The simulations started with RMSD around 8 Å. It took about 1 μ s before folded structures (with RMSD below 2 Å) were seen. This time scale is reasonable for a long helix [50]. The first trajectory (shown as a black line in Fig. 7) showed an approximately three-stage folding processes. Together with the dictionary of secondary structure of proteins (DSSP) [51] analysis (see the top panel in Fig. 8), we found that two helices (residues 3–6 and 10–20) formed very quickly and remained very stable and unmerged in the first 400 ns. The short helix was mainly a 3_{10} helix. The long helix grew towards the C-terminus, and elongated by eight residues in the next 100 ns. This structure held for about 400 ns. Finally, in the last 100 ns, these two helices merged together and the RMSD dropped to below 2 Å. Another trajectory showed only some trivial differences. In the very beginning of this simulation, three helices (residues 3–7, 10–20, 23–28) formed quickly. The helix close to the N-terminus is an α -helix, while the other two are 3_{10} helices. The helix close to the C-terminus (residues 23–28) unfolded at 370 ns, and remained as a random coil in the next 500 ns. At the same time, the central helix switched to an α -helix, while the first helix changed into an 3_{10} helix. At 700 ns, these two helices merged together, and extended towards the C-terminus. The folded structure was arrived at after 900 ns.

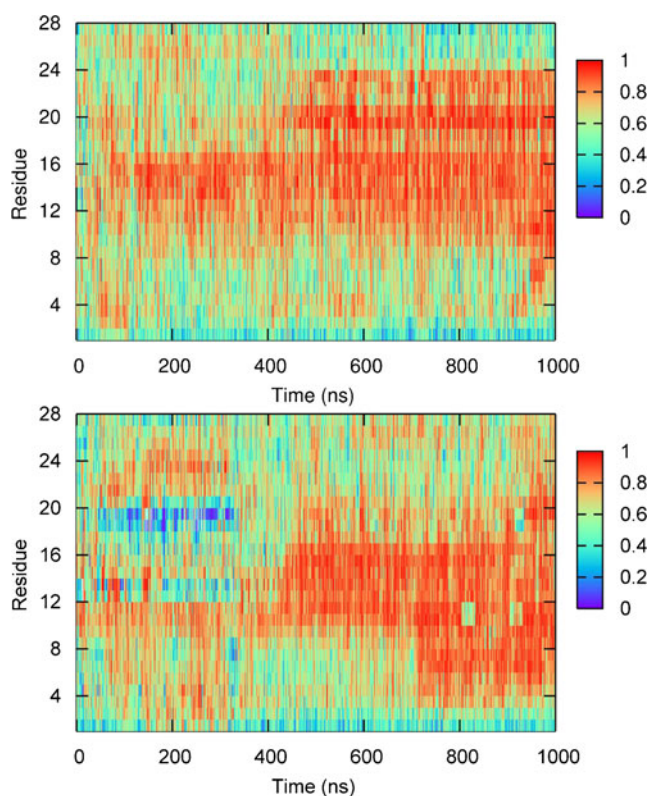


Fig. 9 Variations of local TFE concentration (LTC) in the folding simulations

The variation of the LTC around the peptide is shown in Fig. 9. At the beginnings of both simulations, the LTC was around 50 %, but it gradually increased during the simulations. We noticed that the increase in LTC and the formation of helical residues were highly correlated. In the final stage of these simulations, the LTC approached 100 % for the central residues, and decayed toward two terminals. This result was consistent with the equilibrium MD simulation. This observation serves as direct evidence that TFE is an effective inducer of helix.

Conclusions

In this work, both equilibrium and folding simulations were carried out to study the solvation effect of TFE on the structure of a small peptide. The structure of this peptide, which is a long helix, was solved in 50 % (v/v) TFE. In equilibrium simulations, both explicit and implicit solvation models have been utilized and compared. Explicit solvent models provide arguably the most accurate description of the solvation effect. The potential function of TFE was produced following the routine of the general AMBER force field. The peptide is partly unfolded in water but adopts an ordered helical conformation in a water/TFE mixture, which is in agreement with the experimental observations [28]. Furthermore, the accumulation of TFE molecules in the immediate vicinity of this peptide was observed, which prevents the formation of alternative hydrogen bonds between water and the peptide. Besides, it provides a low dielectric environment that stabilizes the main chain hydrogen bonds in the peptide.

The folding simulations were performed in explicit water/TFE solvent. These two trajectories show only trivial differences. The folding takes place mainly from central residues, and elongates towards the terminals. The helical segments merge together in the final stage. DSSP and LTC analysis of the trajectories shows that the folding is accompanied by the aggregation of TFE molecules on the peptide surface. This coating effect promotes the formation of local interactions and, as a consequence, an ordered secondary structure. This observation is consistent with the experimental understanding that TFE is an effective trigger of secondary structure, especially of helical conformation.

These simulations have demonstrated the effectiveness of this water/TFE mixed solvent model. However, GB representation of the water/TFE mixed solvent failed to obtain results in agreement with the explicit solvent model. Although GB methodology is successful in modeling simple homogeneous and invariable dielectric environment [46, 52, 53], the water/TFE mixed solvent is more complex and requires a heterogeneous and alterable description of dielectric constant in order to depict the aggregated TFE molecules. So the GB solvent model is not an appropriate choice for water/TFE mixture so far.

Acknowledgments This work is supported by the National Natural Science Foundation of China (Grant No. 21173082) and the Shanghai Rising-Star Program (Grant No. 11QA1402000). We thank Supercomputer Center of East China Normal University for CPU time support.

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